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## On some constants related to the harmonic zeta function

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**Abstract** We study a range of constants that are closely linked to the harmonic zeta function  $\zeta_H$ . In addition, by means of a new representation of  $\zeta_H$ , we derive possibly new integral formulas for  $\zeta(n)$  and  $\zeta'(n)$  and new evaluations for a variety of integrals of the same type as those previously considered by Glasser, Manna, and Oloa.

**Keywords** Harmonic zeta function; Stieltjes constants; zeta values; Glasser-Manna-Oloa integrals, Ramanujan summation of series.

## 1 Introduction

Let  $\zeta_H$  be the harmonic zeta function defined for Re(s) > 1 by

$$\zeta_H(s) := \sum_{n=1}^{\infty} \frac{H_n}{n^s} \,,$$

where, for all  $n \geq 1$ ,

$$H_n = 1 + \frac{1}{2} + \dots + \frac{1}{n}$$

are the classical harmonic numbers. Four decades ago, Apostol and Vu [1] showed that this function could be continued as a meromorphic function with a double pole at s=1, and an infinite number of simple poles at s=0 and s=1-2k for each integer  $k \geq 1$ . The Laurent expansion of the harmonic zeta function  $\zeta_H$  around its double pole s=1 can be written as

$$\zeta_H(s) = \frac{1}{(s-1)^2} + \frac{\gamma}{s-1} + \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \gamma_H^{(k)}(s-1)^k \quad (0 < |s-1| < 1),$$

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where  $\gamma = -\Gamma'(1)$  is Euler's constant, and the coefficients  $\gamma_H^{(k)}$  are called harmonic Stieltjes constants by analogy with the classical Stieltjes constants  $\gamma_k$ . A common definition of these constants is

$$\gamma_k = \lim_{s \to 1} \left\{ (-1)^k \zeta^{(k)}(s) - \frac{k!}{(s-1)^{k+1}} \right\} \qquad (k \ge 0),$$

where  $\zeta^{(k)}(s)$  is the kth derivative of the Riemann zeta function. In section 3, we give a new evaluation of the harmonic Stieltjes constants  $\gamma_H^{(k)}$  (see Corollary 1) which arises from a new expression of  $\zeta_H(s)$ . From now on and throughout the text, we will use the lighter notation  $\gamma^{(k)}$  in place of  $\gamma_H^{(k)}$ . An explicit expression of  $\gamma^{(0)}$  is given by the following formula [9, Equation (6)]:

$$\gamma^{(0)} = \frac{1}{2}\gamma^2 + \frac{1}{2}\zeta(2) = \frac{1}{2}\Gamma''(1). \tag{1}$$

It is worth noting that this nice identity extends to the case of the height 1 multiple zeta functions  $\zeta(s, 1, ..., 1)$  [17, Corollary 4.1]. Regarding the classical Stieltjes constants, we recall the well-known asymptotic formula

$$\gamma_k = \lim_{N \to \infty} \left\{ \sum_{n=1}^N \frac{(\log n)^k}{n} - \frac{(\log N)^{k+1}}{k+1} \right\} \qquad (k \ge 0).$$

From our point of view, it is important to point out that  $\gamma_k$  is nothing else than the  $\mathcal{R}$ -sum (i.e. the sum of the series in the sense of Ramanujan's summation method) of the divergent series  $\sum_{n\geqslant 1} \frac{(\log n)^k}{n}$  [8, p. 67]. The existence of a similar asymptotic formula for  $\gamma^{(k)}$ :

$$\gamma^{(k)} = \lim_{N \to \infty} \left\{ \sum_{n=1}^{N} \frac{H_n(\log n)^k}{n} - \frac{(\log N)^{k+2}}{k+2} - \gamma \frac{(\log N)^{k+1}}{k+1} \right\} \qquad (k \ge 0),$$

which results from the extension of a formula of Briggs and Buschman [3, 7], strongly suggests that the constant  $\gamma^{(k)}$  is closely linked to the  $\mathcal{R}$ -sum of the divergent series  $\sum_{n\geqslant 1} \frac{H_n (\log n)^k}{n}$ . This is indeed the case: such a relation exists although it is more intricate than in the classical case. We have already encountered in earlier studies [8, Equation (3.23)], [9, Equation (4)] the following relation:

$$r_0 = \gamma^{(0)} + \tau_1 + \gamma_1 - \zeta(2) = \frac{1}{2}\gamma^2 - \frac{1}{2}\zeta(2) + \gamma_1 + \tau_1, \qquad (2)$$

where

$$r_0 := \sum_{n>1}^{\mathcal{R}} \frac{H_n}{n}$$
 and  $\tau_1 = \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\zeta(n+1)}{n} = -\sum_{n=2}^{\infty} \zeta'(n)$ .

We will see in the next section to what extent our formula (2) can be generalized (see Proposition 1). In the subsequent section, the new representation of  $\zeta_H$  provided by Theorem 1, which is an improvement of a formula due to Boyadzhiev, Gadiyar and Padma [6], allows us to derive new integral formulas for  $\zeta(n)$  and  $\zeta'(n)$  (see Corollary 2), as well as new evaluations for a variety of integrals of the same type as those previously considered by Glasser and Manna [14], and Oloa [16] (see Corollaries 4–6).

# 2 How the constants $\gamma^{(k)}$ and $\sum_{n\geqslant 1}^{\mathcal{R}} \frac{H_n (\log n)^k}{n}$ are linked together

The forthcoming Proposition 1 successfully generalizes the previous relation (2).

**Proposition 1.** Let k be a natural number. If  $r_k$  and  $\tau_k$  are defined by

$$r_k := \sum_{n>1}^{\mathcal{R}} \frac{H_n(\log n)^k}{n} \qquad (k \ge 0),$$

and

$$\tau_k := \sum_{n=1}^{\infty} (-1)^{n+k} \frac{\zeta(n+1)}{n^k} \qquad (k \ge 1),$$

then

$$r_0 = \gamma^{(0)} + \tau_1 + \gamma_1 - \zeta(2) ,$$

and the following relations hold true for all positive values of k:

$$r_{2k-1} = \gamma^{(2k-1)} + (2k-1)! \,\tau_{2k} + \frac{\gamma_{2k}}{2k} + \sum_{n=1}^{k} \frac{(2k-1)!}{(2k-2n)!} \times \frac{(2^{2n-1}-1)}{2^{2n-2}} \,\gamma_{2k-2n} \,\zeta(2n) \,, \tag{3}$$

and

$$r_{2k} = \gamma^{(2k)} + (2k)! \, \tau_{2k+1} + \frac{\gamma_{2k+1}}{2k+1} + \sum_{n=1}^{k} \frac{(2k)!}{(2k-2n+1)!} \times \frac{(2^{2n-1}-1)}{2^{2n-2}} \, \gamma_{2k-2n+1} \zeta(2n) - (2k)! \frac{(2^{2k+1}-1)}{2^{2k}} \zeta(2k+2) \,. \tag{4}$$

In particular,

$$r_{1} = \gamma^{(1)} + \tau_{2} + \frac{1}{2}\gamma_{2} + \gamma\zeta(2),$$

$$r_{2} = \gamma^{(2)} + 2\tau_{3} + \frac{1}{3}\gamma_{3} + 2\gamma_{1}\zeta(2) - \frac{7}{2}\zeta(4),$$

$$r_{3} = \gamma^{(3)} + 6\tau_{4} + \frac{1}{4}\gamma_{4} + 3\gamma_{2}\zeta(2) + \frac{21}{2}\gamma\zeta(4),$$

$$r_{4} = \gamma^{(4)} + 24\tau_{5} + \frac{1}{5}\gamma_{5} + 4\gamma_{3}\zeta(2) + 42\gamma_{1}\zeta(4) - \frac{93}{2}\zeta(6).$$

**Example 1.** Numerical evaluations of  $\gamma^{(k)}$ ,  $r_k$  and  $\tau_k$  for  $k \leq 4$  are given below.

$$\begin{array}{llll} \gamma^{(0)} = 0.989055\ldots & r_0 = 0.529052\ldots & \tau_0 \text{ is undefined} \\ \gamma^{(1)} = 0.400761\ldots & r_1 = -0.078850\ldots & \tau_1 = 1.257746\ldots \\ \gamma^{(2)} = 0.971304\ldots & r_2 = -0.008095\ldots & \tau_2 = -1.424248\ldots \\ \gamma^{(3)} = 2.972584\ldots & r_3 = 0.003653\ldots & \tau_3 = 1.523800\ldots \\ \gamma^{(4)} = 11.940501\ldots & r_4 = 0.002816\ldots & \tau_4 = -1.580233\ldots \end{array}$$

**Remark 1.** Although not used afterwards, we also mention another expression of the constant  $\tau_k$  which can be easily deduced from its definition:

$$\tau_k = (-1)^k \sum_{n=1}^{\infty} \frac{1}{n} \operatorname{Li}_k \left( -\frac{1}{n} \right) \qquad (k \ge 1),$$

where  $\text{Li}_k$  denotes the kth polylogarithm. In particular, for k=1, the formula

$$\tau_1 = \sum_{n=1}^{\infty} \frac{1}{n} \log \left( 1 + \frac{1}{n} \right)$$

is regained. Note that the definition of  $\tau_k$  given here differs from that given in [9] when k > 1.

Proof of Proposition 1. The key formula to derive the general relations (3) and (4) is the decomposition of  $\sum_{n\geqslant 1}^{\mathcal{R}} \frac{H_n}{n^s}$  given by [9, Theorem 1] which is reproduced below:

$$\sum_{n\geq 1}^{\mathcal{R}} \frac{H_n}{n^s} = \frac{\pi}{\sin(\pi s)} \, \zeta(s) + \int_0^1 \frac{\psi(x+1) + \gamma}{x^s} \, dx + \zeta_H(s) \,, \tag{5}$$

where  $\psi(x) = \frac{d}{dx} \log \Gamma(x)$  is the digamma function. This formula applies to all complex numbers s such that Re(s) < 2 and  $s \neq 1, 0, 1 - 2k$  for each integer k > 0. Fortunately, the Laurent expansion of each component in (5) can be written explicitly.

a) The expansion of  $\frac{\pi}{\sin(\pi s)} \zeta(s)$  at s=1 can be obtained as follows: first, we write the successive equations:

$$\frac{-\pi}{\sin(\pi s)} = \frac{\pi}{\sin(\pi(s-1))} = \frac{\exp(i\pi(s-1))}{s-1} \frac{2i\pi(s-1)}{\exp(2i\pi(s-1))-1}$$
$$= \frac{1}{s-1} \left( \sum_{k\geq 0} i^k \pi^k \frac{1}{k!} (s-1)^k \right) \left( \sum_{k\geq 0} i^k (2\pi)^k \frac{B_k}{k!} (s-1)^k \right)$$
$$= \frac{1}{s-1} + \sum_{k\geq 1} (-1)^k \left( \sum_{j=0}^{2k} {2k \choose j} 2^j B_j \right) \frac{\pi^{2k}}{(2k)!} (s-1)^{2k-1},$$

where  $B_j$  are the Bernoulli numbers. Euler's identity:

$$\zeta(2k) = (-1)^{k+1} 2^{2k-1} B_{2k} \frac{\pi^{2k}}{(2k)!} \qquad (k \ge 1),$$

then allows us to rewrite this expansion as follows:

$$\frac{-\pi}{\sin(\pi s)} = \frac{1}{s-1} - \sum_{k>1} \frac{2^{1-2k}}{B_{2k}} \sum_{j=0}^{2k} {2k \choose j} 2^j B_j \zeta(2k)(s-1)^{2k-1}.$$

Morover, the latter expression can be simplified thanks to the identity

$$\sum_{j=0}^{k} {k \choose j} 2^{j} B_{j} = 2^{k} B_{k} (\frac{1}{2}) = 2(1 - 2^{k-1}) B_{k} \qquad (k \ge 2).$$

Hence, the expansion of  $\frac{\pi}{\sin(\pi s)}$  at s=1 is given by

$$\frac{\pi}{\sin(\pi s)} = -\frac{1}{s-1} - \sum_{k=1}^{\infty} \frac{2^{2k-1} - 1}{2^{2k-2}} \zeta(2k)(s-1)^{2k-1}.$$

On the other hand, the expansion of  $\zeta(s)$  at s=1 is

$$\zeta(s) = \frac{1}{s-1} + \gamma + \sum_{k=1}^{\infty} \frac{(-1)^k}{k!} \gamma_k (s-1)^k,$$

where  $\gamma_k$  are the Stieltjes constants. The expansion of  $\frac{\pi}{\sin(\pi s)} \zeta(s)$  is then obtained by Cauchy product. We have

$$\frac{\pi}{\sin(\pi s)} \zeta(s) = -\frac{1}{(s-1)^2} - \frac{\gamma}{(s-1)} + \gamma_1 - \zeta(2) 
- \left(\frac{1}{2}\gamma_2 + \gamma\zeta(2)\right) (s-1) + \left(\frac{1}{6}\gamma_3 + \gamma_1\zeta(2) - \frac{7}{4}\zeta(4)\right) (s-1)^2 
- \left(\frac{1}{24}\gamma_4 + \frac{1}{2}\gamma_2\zeta(2) + \frac{7}{4}\gamma\zeta(4)\right) (s-1)^3 
+ \left(\frac{1}{120}\gamma_5 + \frac{1}{6}\gamma_3\zeta(2) + \frac{7}{4}\gamma_1\zeta(4) - \frac{31}{16}\zeta(6)\right) (s-1)^4 - \cdots (6)$$

In (6), the coefficient of  $(s-1)^{2k}$  and  $(s-1)^{2k-1}$  are respectively

$$\frac{\gamma_{2k+1}}{(2k+1)!} + \sum_{j=1}^{k} \frac{\gamma_{2k-2j+1}}{(2k-2j+1)!} \times \frac{(2^{2j-1}-1)}{2^{2j-2}} \zeta(2j) - \frac{(2^{2k+1}-1)}{2^{2k}} \zeta(2k+2),$$

and

$$-\frac{\gamma_{2k}}{(2k)!} - \sum_{j=1}^{k} \frac{\gamma_{2k-2j}}{(2k-2j)!} \times \frac{(2^{2j-1}-1)}{2^{2j-2}} \zeta(2j).$$

b) The well-known Taylor series expansion of the digamma function:

$$\psi(x+1) + \gamma = \sum_{n=1}^{\infty} (-1)^{n+1} \zeta(n+1) x^n \qquad (|x| < 1),$$

allows us write

$$\int_0^1 \frac{\psi(x+1) + \gamma}{x^s} dx = \sum_{n=1}^\infty (-1)^{n+1} \frac{\zeta(n+1)}{n - (s-1)}.$$

This leads to the following expansion:

$$\int_0^1 \frac{\psi(x+1) + \gamma}{x^s} dx = \sum_{k=0}^\infty \left( \sum_{n=1}^\infty (-1)^{n+1} \frac{\zeta(n+1)}{n^{k+1}} \right) (s-1)^k$$
$$= \sum_{k=0}^\infty \frac{(-1)^k}{k!} \times k! \, \tau_{k+1} \, (s-1)^k \qquad (|s-1| < 1). \quad (7)$$

c) To obtain the expansion of  $\sum_{n\geqslant 1}^{\mathcal{R}}\frac{H_n}{n^s}$  at s=1, we proceed as follows: first we write the identities

$$\frac{H_n}{n^s} = \frac{H_n}{n} e^{-(s-1)\log n} = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \times \frac{H_n(\log n)^k}{n} (s-1)^k,$$

so that

$$\sum_{n\geqslant 1}^{\mathcal{R}} \frac{H_n}{n^s} = \sum_{n>1}^{\mathcal{R}} \left( \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \times \frac{H_n(\log n)^k}{n} \left( s - 1 \right)^k \right) .$$

Moreover, [8, Theorem 9] allows us interchange  $\sum_{n\geqslant 1}^{\mathcal{R}}$  and  $\sum_{k=0}^{\infty}$ . This leads to the expansion:

$$\sum_{n\geqslant 1}^{\mathcal{R}} \frac{H_n}{n^s} = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \left( \sum_{n\geq 1}^{\mathcal{R}} \frac{H_n(\log n)^k}{n} \right) (s-1)^k = \sum_{k=0}^{\infty} \frac{(-1)^k}{k!} r_k (s-1)^k.$$
 (8)

The desired formulas (3) and (4) can then be easily obtained by combining the previous expansions (6), (7), and (8).

## 3 A new expresssion of $\zeta_H(s)$ and its implications

In this section, we provide a new expression of the harmonic zeta function (Theorem 1) which allows us to give a new evaluation of the harmonic Stieltjes constants (Corollary 1). Thanks to this representation of  $\zeta_H$ , we also derive new integral formulas for  $\zeta(n)$  and  $\zeta'(n)$  (Corollary 2), and new evaluations for Glasser-Manna-Oloa type integrals (Corollaries 4–6).

**Theorem 1.** For all complex numbers s in  $\mathbb{C} \setminus \mathbb{Z}$ , we have

$$\zeta_H(s) = \pi \cot(\pi s) \, \zeta(s) + \zeta(s+1) + \Gamma(1-s)\Phi(s) \tag{9}$$

with

$$\Phi(s) := -\frac{1}{2i\pi} \int_{-\pi}^{\pi} x \left( \text{Log}(1 + e^{ix}) \right)^{s-1} dx,$$
 (10)

where Log denotes the principal value of the complex logarithm.

*Proof.* Let

$$z = \operatorname{Log}(1 + e^{ix}) = \log\left(2\cos\left(\frac{x}{2}\right)\right) + i\frac{x}{2}, \quad -\pi < x < \pi.$$

When x varies from  $-\pi$  to  $\pi$ , the variable z traverses the path L defined by the parametric equations  $\text{Re } z = \log((2\cos(x/2)))$  and Im z = x/2. This path extends from  $-\infty$  below the line Im z = 0, passes through the point  $(\log 2, 0)$ , then extends back to  $-\infty$  above the line Im z = 0. Since L is homotopic to the Hankel contour C defined in [5, 6], then by [6, Equation (25)], we have, for  $s \in \mathbb{C} \setminus \mathbb{Z}$ , the representation

$$\zeta_H(s) = \pi \cot(\pi s) \, \zeta(s) + \zeta(s+1) + \Gamma(1-s)\phi(s) - \psi(1-s)\zeta(s) - \zeta'(s) \,, \tag{11}$$

where

$$\phi(s) := \frac{1}{2i\pi} \int_{L} \frac{z^{s-1}e^{z}}{e^{z}-1} \operatorname{Log}\left(\frac{e^{z}-1}{z}\right) dz.$$

We now show that equations (9) and (11) are equivalent. To do this, we proceed as follows. Differentiation of the integral representation

$$\zeta(s) = \frac{\Gamma(1-s)}{2i\pi} \int_{L} \frac{z^{s-1}e^{z}}{1-e^{z}} dz$$

[5, Equation (2.4)], leads to the following identity:

$$\psi(1-s)\zeta(s) + \zeta'(s) = \frac{\Gamma(1-s)}{2i\pi} \int_{L} \frac{z^{s-1}e^{z} \operatorname{Log}(z)}{1-e^{z}} dz \qquad (s \neq 1, 2, 3, \ldots)$$
 (12)

[5, Equation (2.7)]. Splitting the right-hand member of (12), we can write

$$\frac{1}{2i\pi} \int_{L} \frac{z^{s-1} e^{z} \operatorname{Log}(z)}{1 - e^{z}} dz = \phi(s) - \Phi(s), \qquad (13)$$

with

$$\Phi(s) = \frac{1}{2i\pi} \int_{L} \frac{z^{s-1}e^{z}}{e^{z} - 1} \text{Log}(e^{z} - 1) dz.$$

From (12) and (13), we deduce the identity

$$\Gamma(1-s)\phi(s) - \psi(1-s)\zeta(s) - \zeta'(s) = \Gamma(1-s)\Phi(s).$$

Moreover, the substitution

$$z = \text{Log}(1 + e^{ix}), \quad -\pi < x < \pi,$$

in the contour integral defining  $\Phi$  gives this function the equivalent form (10), proving that (9) and (11) are equivalent equations.

## 3.1 New formulas for $\gamma^{(k)}$

We deduce from (9) an evaluation of  $\gamma^{(k)}$  which seems to us more meaningful than the one given in [15, Equation (21)].

Corollary 1. For any non-negative integer n, let  $\mathcal{L}_n$  and  $\xi_n$  be defined as follows:

$$\mathcal{L}_n := \Phi^{(n)}(1) = -\frac{1}{2i\pi} \int_{-\pi}^{\pi} x \operatorname{Log}^n(\operatorname{Log}(1 + e^{ix})) dx,$$

and

$$\xi_n := (-1)^n \Gamma^{(n)}(1) = (-1)^n \int_0^{+\infty} e^{-x} \log^n(x) dx$$
.

Then we have

$$\gamma^{(0)} = -\mathcal{L}_1 - \gamma_1 - \zeta(2), \qquad (14)$$

and, for all positive integers k, the following general relations hold true:

$$\gamma^{(2k-1)} = -\zeta^{(2k-1)}(2) - \frac{\gamma_{2k}}{2k} + \frac{1}{2k} \sum_{n=0}^{2k-1} {2k \choose n} \xi_n \mathcal{L}_{2k-n} + \sum_{n=1}^k \frac{2(2k-1)!}{(2k-2n)!} \gamma_{2k-2n} \zeta(2n), \quad (15)$$

and

$$\gamma^{(2k)} = \zeta^{(2k)}(2) - \frac{\gamma_{2k+1}}{2k+1} - \frac{1}{2k+1} \sum_{n=0}^{2k} {2k+1 \choose n} \xi_n \mathcal{L}_{2k+1-n} + \sum_{n=1}^{k} \frac{2(2k)!}{(2k+1-2n)!} \gamma_{2k+1-2n} \zeta(2n) - 2(2k)! \zeta(2k+2). \quad (16)$$

Proof of Corollary 1. The key formula to derive the general relations (15) and (16) is the representation of  $\zeta_H(s)$  given by (9). Fortunately, the Laurent series expansion of each component in (9) can be written explicitly.

a) The Laurent expansions of  $\pi \cot(\pi s)$  and  $\zeta(s)$  at s=1 are respectively

$$\pi \cot(\pi s) = \frac{1}{s-1} - \sum_{k>1} 2\zeta(2k)(s-1)^{2k-1},$$

and

$$\zeta(s) = \frac{1}{s-1} + \gamma + \sum_{k>1} \frac{(-1)^k}{k!} \gamma_k (s-1)^k,$$

where the coefficients  $\gamma_k$  are the Stieltjes constants. The expansion of  $\pi \cot(\pi s) \zeta(s)$  is then obtained by Cauchy product:

$$\pi \cot(\pi s) \zeta(s) = \frac{1}{(s-1)^2} + \frac{\gamma}{s-1} - \gamma_1 - 2\zeta(2) - \left(-\frac{1}{2}\gamma_2 + 2\gamma\zeta(2)\right)(s-1) + \frac{1}{2}\left(-\frac{1}{3}\gamma_3 + 4\zeta(2)\gamma_1 - 4\zeta(4)\right)(s-1)^2 - \frac{1}{6}\left(-\frac{1}{4}\gamma_4 + 6\zeta(2)\gamma_2 + 12\gamma\zeta(4)\right)(s-1)^3 + \cdots$$
 (17)

b) It follows from the definition of  $\xi_k$  as  $(-1)^k \Gamma^{(k)}(1)$  that the Laurent expansion of  $\Gamma(1-s)$  at s=1 is given by

$$\Gamma(1-s) = -\frac{1}{s-1} - \gamma - \sum_{k=2}^{\infty} \frac{\xi_k}{k!} (s-1)^{k-1} \qquad (|s-1| < 1),$$

On the other hand, the function  $\Phi$  defined by (10) is an entire function of s with  $\Phi(1) = \mathcal{L}_0 = 0$ , and the definition of  $\mathcal{L}_k$  as  $\Phi^{(k)}(1)$  implies that the Taylor expansion of  $\Phi(s)$  at s = 1 is given by

$$\Phi(s) = \sum_{k=1}^{\infty} \frac{\mathcal{L}_k}{k!} (s-1)^k.$$

The Laurent expansion of  $\Gamma(1-s)\Phi(s)$  then follows by Cauchy product:

$$\Gamma(1-s)\Phi(s) = -\mathcal{L}_1 - \left(\frac{1}{2}\mathcal{L}_2 + \gamma\mathcal{L}_1\right)(s-1) + \frac{1}{2}\left(-\frac{1}{3}\mathcal{L}_3 - \gamma\mathcal{L}_2 - 2\gamma^{(0)}\mathcal{L}_1\right)(s-1)^2 - \frac{1}{6}\left(\frac{1}{4}\mathcal{L}_4 + \gamma\mathcal{L}_3 + 3\gamma^{(0)}\mathcal{L}_2 + (2\zeta(3) + 3\gamma\zeta(2) + \gamma^3)\mathcal{L}_1\right)(s-1)^3 + \cdots$$
(18)

c) The Taylor expansion of  $\zeta(s+1)$  at s=1 can be written as follows:

$$\zeta(s+1) = \zeta(2) + \sum_{k=1}^{\infty} \frac{(-1)^k}{k!} (-1)^k \zeta^{(k)}(2) (s-1)^k.$$
 (19)

By assembling equations (17), (18), and (19) above, we obtain, by identifying the constant term, the relation (14). In the same way, the general formulae (15) and (16) are derived by identifying the coefficients of higher degree.  $\Box$ 

**Remark 2.** a) Combining (14) with the expression of  $\gamma^{(0)}$  given by (1), we deduce the explicit evaluation:

$$\mathcal{L}_1 = -\frac{1}{2i\pi} \int_{-\pi}^{\pi} x \operatorname{Log}(\operatorname{Log}(1 + e^{ix})) dx = -\frac{3}{2}\zeta(2) - \frac{1}{2}\gamma^2 - \gamma_1.$$
 (20)

b) The constants  $\xi_n$  involved in formulae (15)–(16) have a closed form expression in terms of Euler's constant  $\gamma$  and  $\zeta(2), \ldots, \zeta(n)$  [2, p. 227], [10, p. 12]. More precisely, we have  $\xi_0 = 1$ ,  $\xi_1 = \gamma$ , and the recursion formula

$$\xi_{n+1} = \gamma \xi_n + \sum_{k=1}^n \frac{n!}{(n-k)!} \zeta(k+1) \xi_{n-k} \qquad (n \ge 1).$$

so that

$$\xi_2 = \gamma^2 + \zeta(2) = 2\gamma^{(0)},$$
  
 $\xi_3 = \gamma^3 + 3\gamma\zeta(2) + 2\zeta(3),$   
 $\xi_4 = \gamma^4 + 6\gamma^2\zeta(2) + 8\gamma\zeta(3) + 3(\zeta(2))^2 + 6\zeta(4), \text{ etc.}$ 

A clever reformulation of this result is

$$\xi_n = P_n(\gamma, \zeta(2), \dots, \zeta(n)) \qquad (n \ge 2),$$

where the polynomials  $P_n$  are defined by the exponential generating function

$$\exp\left(\sum_{k=1}^{\infty} x_k \frac{t^k}{k}\right) = \sum_{n=0}^{\infty} P_k(x_1, x_2, \dots, x_n) \frac{t^n}{n!}.$$

c) By mean of a theorem of Ramanujan [11, Theorem 4] whose proof is given in [4, p. 224], we can write the following identity:

$$(-1)^n \zeta^{(n)}(2) = n! + \gamma_n + \sum_{k=1}^{\infty} \frac{(-1)^k}{k!} \gamma_{k+n} \qquad (n \ge 1).$$

**Example 2.** Applying formulas (15) and (16) to the case k = 1, and using (20), we obtain respectively

$$\gamma^{(1)} = -\zeta'(2) - \frac{1}{2}\gamma^3 - \frac{1}{2}\gamma_2 + \frac{1}{2}\mathcal{L}_2 + \frac{1}{2}\gamma\zeta(2) - \gamma\gamma_1,$$

and

$$\gamma^{(2)} = \zeta''(2) + \frac{1}{2}\gamma^4 - \frac{1}{3}\gamma_3 - \frac{1}{3}\mathcal{L}_3 - \gamma\mathcal{L}_2 + (2\gamma^2 + 5\gamma_1)\zeta(2) + \gamma^2\gamma_1 - \frac{1}{10}(\zeta(2))^2.$$

Numerical evaluations of the constants  $\mathcal{L}_2$  and  $\mathcal{L}_3$  are

$$\mathcal{L}_2 = -1.924491\dots$$
 and  $\mathcal{L}_3 = 7.158075\dots$ 

### **3.2** New formulas for $\zeta(n)$ and $\zeta'(n)$

Corollary 2. For all integers  $n \geq 2$ , we have

$$\zeta(n) = \frac{(-1)^n}{(n-1)!} \times \frac{1}{2i\pi} \int_{-\pi}^{\pi} x \left( \text{Log}(1+e^{ix}) \right)^{n-1} dx, \qquad (21)$$

and

$$\zeta'(n) = \frac{(-1)^n}{(n-1)!} \times \frac{1}{2i\pi} \int_{-\pi}^{\pi} x \operatorname{Log}\left(\operatorname{Log}(1+e^{ix})\right) \left(\operatorname{Log}(1+e^{ix})\right)^{n-1} dx + \zeta_H(n) - \zeta(n) \psi(n) - \zeta(n+1).$$
 (22)

Proof of Corollary 2. For  $n \geq 2$ , we can write

$$\pi \cot(\pi s) \zeta(s) = \frac{\zeta(n)}{s-n} + \zeta'(n) + O(s-n),$$

$$\Gamma(1-s) = \frac{(-1)^n}{(n-1)!} \left(\frac{1}{s-n} + (\gamma - H_{n-1})\right) + O(s-n),$$

and

$$\Gamma(1-s)\,\Phi(s) = \frac{(-1)^n}{(n-1)!}\frac{\Phi(n)}{s-n} + \frac{(-1)^n}{(n-1)!}\left(\Phi'(n) + \Phi(n)(\gamma - H_{n-1})\right) + \mathcal{O}(s-n)\,.$$

By applying formula (9) around s = n, we deduce the equation

$$\zeta_H(n) + \mathcal{O}(s-n) = \left(\zeta(n) + \frac{(-1)^n}{(n-1)!}\Phi(n)\right) \frac{1}{s-n} + \frac{(-1)^n}{(n-1)!} \left(\Phi'(n) + \Phi(n)(\gamma - H_{n-1})\right) + \zeta'(n) + \zeta(n+1) + \mathcal{O}(s-n).$$

This leads to the identities

$$\zeta(n) + \frac{(-1)^n}{(n-1)!}\Phi(n) = 0$$
,

and

$$\zeta_H(n) = \frac{(-1)^n}{(n-1)!} \left( \Phi'(n) + \Phi(n)(\gamma - H_{n-1}) \right) + \zeta'(n) + \zeta(n+1) ,$$

thus proving formulas (21) and (22).

Corollary 3. For all positive integers n, we have

$$\Phi(2n+1) = (2n)! \, \zeta(2n+1) \,, \tag{23}$$

and

$$-\Phi'(2n) = (2n-1)! \left(\zeta'(2n) + \zeta(2n)\psi(2n) - n\zeta(2n+1) + \sum_{k=1}^{n-1} \zeta(2n-k)\zeta(k+1)\right). \tag{24}$$

Proof of Corollary 3. By Euler's formula [4, p. 252], we have

$$\zeta_H(2n) = (n+1)\zeta(2n+1) - \sum_{k=1}^{n-1} \zeta(2n-k)\zeta(k+1),$$

and thus formulas (23) and (24) derive respectively from (21) and (22).

**Example 3.** Applying (23) and (24) for n = 1, we obtain

$$\Phi(3) = 2\zeta(3) = \zeta_H(2) = -\frac{1}{2i\pi} \int_{-\pi}^{\pi} x \operatorname{Log}^2(1 + e^{ix}) dx$$

and

$$-\Phi'(2) = \zeta'(2) + \zeta(2)(1-\gamma) - \zeta(3) = \frac{1}{2i\pi} \int_{-\pi}^{\pi} x \operatorname{Log}(1+e^{ix}) \operatorname{Log}\left(\operatorname{Log}(1+e^{ix})\right) dx.$$

**Remark 3.** It should be noted that our formula (21) for  $\zeta(n)$  is similar to but different from that given in [5, Equation (2.6)] which, after substitution  $z = \text{Log}(1 + e^{ix})$ , translates into

$$\zeta(n) = \frac{(-1)^{n-1}}{(n-1)!} \times \frac{1}{2\pi} \int_{-\pi}^{\pi} \text{Log}\left(\text{Log}(1+e^{ix})\right) \left(\text{Log}(1+e^{ix})\right)^{n-1} dx.$$

#### 3.3 New formulas for Glasser-Manna-Oloa type integrals

Using the relation between  $\Phi(s)$  and  $\zeta_H(s)$  arising from formula (9) in a neighborhood of s = -n, we derive new formulas for Glasser-Manna-Oloa type integrals.

Corollary 4. We have

$$\Phi(0) = -\frac{1}{2i\pi} \int_{-\pi}^{\pi} \frac{x}{\text{Log}(1+e^{ix})} dx = \frac{1}{2} \log(2\pi) - \frac{1}{2}\gamma + \frac{1}{2},$$
 (25)

and

$$\Phi'(0) = -\frac{1}{2i\pi} \int_{-\pi}^{\pi} \frac{x \operatorname{Log}(\operatorname{Log}(1 + e^{ix}))}{\operatorname{Log}(1 + e^{ix})} dx 
= \frac{1}{2}\gamma_1 + \frac{1}{4}\operatorname{log}^2(2\pi) + \frac{1}{4}\gamma^2 - \frac{1}{2}\gamma - \frac{1}{2}\gamma \operatorname{log}(2\pi) - \frac{7}{8}\zeta(2) + \beta, \qquad (26)$$

where  $\beta$  is the linear coefficient in the Laurent expansion of  $\zeta_H$  at s=0.

*Proof of Corollary* 4. Around 0, we have the decomposition given by (9):

$$\zeta_H(s) = \zeta(s+1) + \pi \cot(\pi s) \zeta(s) + \Gamma(1-s)\Phi(s),$$

and the expansion [9, Equation (8)]

$$\zeta_H(s) = \frac{1}{2s} + \frac{1}{2}\gamma + \frac{1}{2} + \beta s + O(s^2).$$

On the other side, we have the expansions

$$\zeta(s+1) = \frac{1}{s} + \gamma - \gamma_1 s + \mathcal{O}(s^2),$$

$$\pi \cot(\pi s) \zeta(s) = -\frac{1}{2s} - \frac{1}{2} \log(2\pi) + \left(\frac{1}{2}\gamma_1 - \frac{1}{4} \log^2(2\pi) + \frac{1}{4}\gamma^2 + \frac{7}{8}\zeta(2)\right) s + \mathcal{O}(s^2),$$

$$\Gamma(1-s)\Phi(s) = \Phi(0) + (\Phi'(0) + \gamma\Phi(0)) + \mathcal{O}(s^2).$$

By identifying the constant coefficient in the expansion of the right-hand member of (9), we deduce the equation

$$\frac{1}{2}\gamma + \frac{1}{2} = \gamma - \frac{1}{2}\log(2\pi) + \Phi(0),$$

which is equivalent to (25). In the same way, formula (26) is derived by identifying the linear coefficient in the expansion of the right-hand member of (9).

**Remark 4.** We derive from [17, Equation (53)] the following evaluation of the constant  $\beta$  which appears in (26):

$$\beta = -\gamma_1^{[2]}(0) - \gamma_1 = 1 + \frac{1}{2}\gamma - \frac{1}{4}\gamma^2 - \gamma_1 + \frac{1}{4}\zeta(2) - \sum_{n=2}^{\infty} \frac{|b_n|}{(n-1)^2} = 1.589935...,$$

where  $b_n$  are the Bernoulli numbers of the second kind defined by way of their generating function

$$\frac{x}{\log(x+1)} = \sum_{n=0}^{\infty} b_n x^n = 1 + \frac{x}{2} - \frac{x^2}{12} + \frac{x^3}{24} - \frac{19x^4}{720} + \frac{3x^5}{160} - \frac{863x^6}{60480} + \cdots$$

**Remark 5.** Formula (25) appears in a slightly different but equivalent form in [14, Equation (3)], due to the identities

$$\Phi(0) = \frac{1}{\pi} \int_{-\pi}^{\pi} \frac{iy \, dy}{2 \log(2 \cos(\frac{y}{2})) + iy}$$

$$= \frac{2}{\pi} \int_{0}^{\pi} \frac{y^{2}}{y^{2} + 4 \log^{2}(2 \cos(\frac{y}{2}))} \, dy$$

$$= \frac{4}{\pi} \int_{0}^{\frac{\pi}{2}} \frac{x^{2}}{x^{2} + \log^{2}(2 \cos(x))} \, dx.$$

Corollary 5. We have

$$\Phi(-1) = -\frac{1}{2i\pi} \int_{-\pi}^{\pi} \frac{x}{\log^2(1 + e^{ix})} dx = \log(\mathcal{A}) - \frac{1}{12}\gamma + \frac{7}{24},$$
 (27)

where  $\mathcal{A} = \exp(\frac{1}{12} - \zeta'(-1))$  is the Glaisher-Kinkelin constant, and for all integers  $n \geq 2$ ,

$$(2n-1)! \Phi(1-2n) = \log(\mathcal{A}_n) - \frac{B_{2n}}{2n} \gamma + \frac{1}{2n} \sum_{k=0}^{2n-2} {2n \choose k} \frac{B_k B_{2n-k}}{2n-k}, \qquad (28)$$

with 
$$\log(\mathcal{A}_n) := \frac{H_{2n-1}B_{2n}}{2n} - \zeta'(1-2n).^{-1}$$

Proof of Corollary 5. Around s = -1, we have the expansion [9, Equation (9)]

$$\zeta_H(s) = -\frac{1}{12(s+1)} - \frac{1}{12}\gamma - \frac{1}{8} + O(s+1).$$

<sup>&</sup>lt;sup>1</sup>The numbers  $A_n$  are called generalized Glaisher-Kinkelin constants [12].

On the other side, we have

$$\zeta(s+1) = -\frac{1}{2} + O(s+1),$$

$$\pi \cot(\pi s) \zeta(s) = \zeta'(-1) + O(s+1),$$

$$\Gamma(1-s)\Phi(s) = \Phi(-1) + O(s+1).$$

Formula (27) is then deduced by identifying the constant coefficient in the expansion of the right-hand member of (9). In the same way, from [9, Proposition 10], we have for  $n \geq 2$  the Laurent expansion around s = 1 - 2n:

$$\zeta_H(s) = \frac{\zeta(1-2n)}{s+2n-1} - \frac{B_{2n}}{2n}\gamma + \frac{H_{2n-1}B_{2n}}{2n} + \frac{1}{2n}\sum_{k=0}^{2n-2} {2n \choose k} \frac{B_k B_{2n-k}}{2n-k} + O(s+2n-1)$$

which, by the same method, allows us to deduce formula (28).

Corollary 6. For all positive integers n, we have the identity

$$\Phi(-2n) = \frac{(-1)^{n+1}}{2} (2\pi)^{-2n} \zeta(2n+1) + \frac{2n+1}{4n} \frac{B_{2n}}{(2n)!}, \qquad (29)$$

from which can be deduced the following relation:

$$2(2n)! \Phi(-2n) = (-1)^{n+1} (2\pi)^{-2n} \Phi(2n+1) + \frac{2n+1}{2n} B_{2n}.$$
 (30)

Proof of Corollary 6. Around s = -2n, we have

$$\zeta_H(s) = -\frac{B_{2n}}{4n} + \frac{B_{2n}}{2} + O(s+2n).$$

On the other side, we have the expansions

$$\zeta(s+1) = -\frac{B_{2n}}{2n} + O(s+2n),$$

$$\pi \cot(\pi s) \zeta(s) = \zeta'(-2n) + O(s+2n),$$

$$\Gamma(1-s)\Phi(s) = (2n)!\Phi(-2n) + O(s+2n).$$

By identifying the constant coefficient in the expansion of the right-hand member of (9), we obtain

$$(2n)! \Phi(-2n) = -\zeta'(-2n) + \frac{2n+1}{4n} B_{2n}$$
$$= \frac{(-1)^{n+1}}{2} (2\pi)^{-2n} (2n)! \zeta(2n+1) + \frac{2n+1}{4n} B_{2n},$$

which is nothing else than (29). Moreover, using the identity

$$\Phi(2n+1) = (2n)! \, \zeta(2n+1)$$

arising from (23) allows us to deduce the relation (30).

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