

## Generalized Glaisher-Kinkelin constants and Blagouchine's integrals

Marc-Antoine Coppo

#### ▶ To cite this version:

Marc-Antoine Coppo. Generalized Glaisher-Kinkelin constants and Blagouchine's integrals. 2021. hal-03197403v11

### HAL Id: hal-03197403 https://univ-cotedazur.hal.science/hal-03197403v11

Preprint submitted on 25 Aug 2021 (v11), last revised 3 Jun 2024 (v23)

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers. L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# Generalized Glaisher-Kinkelin constants and Blagouchine's integrals

Marc-Antoine Coppo<sup>\*</sup>

Université Côte d'Azur, CNRS, LJAD (UMR 7351), Nice, France

Abstract The main purpose of this article is to establish a close connection between two sequences of complex integrals introduced by Blagouchine and some important mathematical constants, namely the Euler-Mascheroni constant, the Cohen-Boyadzhiev constant, and the generalized Glaisher-Kinkelin constants (also known as the Bendersky constants) which occur quite naturally in analysis and number theory. At the end of this study, we also use a formula of Candelpergher to deduce an interesting expression of the alternating series  $\sum_{n\geq 2} (-1)^n \frac{\zeta(n)}{n^s}$  for complex values of s.

**Keywords** Generalized Glaisher-Kinkelin constants, infinite series with zeta values, complex integration.

### 1 Introduction

The main purpose of this article is to highlight the link between the sequence of complex integrals  $\{\mathcal{I}_k\}_{k\geq 0}$  and  $\{\mathcal{J}_k\}_{k\geq 0}$  respectively defined by

$$\mathcal{I}_k = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{3}{2} + ix)}{(2k+1+2ix)\cosh(\pi x)} \, dx$$

and

$$\mathcal{J}_k = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(2k+1+2ix)\cosh(\pi x)} \, dx \,,$$

on one side, and some important mathematical constants, namely the Euler-Mascheroni constant, the Cohen-Boyadzhiev constant, and the generalized Glaisher-Kinkelin constants (also known as the Bendersky constants) which occur quite

<sup>\*</sup>Corresponding author. *Email address:* coppo@unice.fr

naturally in analysis and number theory [7, 10, 11]. To establish these close connections, we make use of a relation (found by Blagouchine [2, Theorem 1]) between the integral  $\mathcal{J}_k$  and the alternating series

$$\nu_k := \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n+k}$$

studied in [8]. Recently, this deep relation has been thoroughly generalized by Candelpergher [6, Eq. (7)], that allows us to give an interesting expression of the alternating series

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n^s}$$

for any complex number s with  $\operatorname{Re}(s) > \frac{1}{2}$  (see formula (11)).

## 2 Generalized Glaisher-Kinkelin constants

**Definition 1** ([1, 10, 11]). For any integer  $k \ge 0$ , the constant  $A_k$  is usually defined by

$$\ln(A_k) = \lim_{N \to \infty} \left\{ \sum_{n=1}^N n^k \ln n - P_k(N) \right\} \,,$$

where  $P_k(N)$  is given by  $P_0(N) = \left(N + \frac{1}{2}\right) \ln N - N$ , and

$$P_k(N) = \left(\frac{N^{k+1}}{k+1} + \frac{N^k}{2} + k! \sum_{j=1}^k \frac{N^{k-j} B_{j+1}}{(j+1)!(k-j)!}\right) \ln N$$
$$- \frac{N^{k+1}}{(k+1)^2} + k! \sum_{j=1}^k \frac{N^{k-j} B_{j+1}}{(j+1)!(k-j)!} \left\{ (1-\delta_{k,j}) \sum_{i=1}^j \frac{1}{k-i+1} \right\} \qquad (k \ge 1),$$

where  $B_j$  is the *j*-th Bernoulli number and  $\delta_{k,j}$  is the Kronecker delta function. The numbers  $A_k$  (for k = 0, 1, 2, ...) are the generalized Glaisher-Kinkelin constants (sometimes called the *Bendersky constants*). In particular, it follows from this definition that

$$\ln(A_0) = \lim_{N \to \infty} \left\{ \sum_{n=1}^N \ln n - \left( N + \frac{1}{2} \right) \ln N + N \right\} \,,$$

and

$$\ln(A_1) = \lim_{N \to \infty} \left\{ \sum_{n=1}^N n \ln n - \left( \frac{N^2}{2} + \frac{N}{2} + \frac{1}{12} \right) \ln N + \frac{N^2}{4} \right\} \,.$$

*Remark* 1. Adamchik [1, Proposition 4] has found a nice expression of theses constants in terms of the derivatives of the Riemann zeta function. More precisely, we have

$$A_k = \exp\left\{\frac{H_k B_{k+1}}{k+1} - \zeta'(-k)\right\} \qquad (k \ge 0),$$
(1)

where  $H_k = \sum_{j=1}^k \frac{1}{j}$  is the *k*-th harmonic number with the usual convention  $H_0 = 0$ . The following relations are easily deduced by differentiation of Riemann's functional equation for the zeta function:

$$\zeta'(-2k) = (-1)^k \frac{(2k)!}{2(2\pi)^{2k}} \zeta(2k+1) \qquad (k \ge 1),$$

and

$$\zeta'(1-2k) = (-1)^{k+1} \frac{(2k)!}{k(2\pi)^{2k}} \zeta'(2k) + \frac{B_{2k}}{2k} \left(H_{2k-1} - \gamma - \ln 2\pi\right) \qquad (k \ge 1).$$

This enable to deduce from Adamchik's formula (1) the expressions

$$A_{2k-1} = \exp\left\{ (-1)^k \frac{(2k)!}{k(2\pi)^{2k}} \zeta'(2k) + \frac{B_{2k}}{2k} \left(\gamma + \ln 2\pi\right) \right\} \qquad (k \ge 1), \qquad (2)$$

and

$$A_{2k} = \exp\left\{ (-1)^{k+1} \frac{(2k)!}{2(2\pi)^{2k}} \zeta(2k+1) \right\} \qquad (k \ge 1).$$
(3)

Example 1. The constant

$$A_0 = \exp(-\zeta'(0)) = \sqrt{2\pi}$$

is the Stirling constant,

$$A_1 = \exp\left(\frac{1}{12} - \zeta'(-1)\right) = \exp\left(-\frac{\zeta'(2)}{2\pi^2} + \frac{\ln(2\pi) + \gamma}{12}\right)$$

is the classical Glaisher-Kinkelin constant, and for k = 2, we have

$$A_2 = \exp(-\zeta'(-2)) = \exp\left(\frac{\zeta(3)}{4\pi^2}\right).$$

Remark 2. Bendersky [3] introduced for the first time the sequence of numbers  $L_k := \ln(A_k)$  without any consideration of their relation with the  $\zeta$ -function. From the point of view of the summation of divergent series, the constants  $L_k$  can be interpreted as follows: if  $\sum_{n\geq 1}^{\mathcal{R}} n^k \ln n$  denotes the  $\mathcal{R}$ -sum of the divergent series  $\sum_{n\geq 1} n^k \ln n$  (i.e. the sum of the series in the sense of Ramanujan's summation method [5]), then, for any integer  $k \geq 0$ , we have

$$\sum_{n\geq 1}^{\mathcal{R}} n^k \ln n = L_k - \frac{H_k B_{k+1}}{k+1} - \frac{1}{(k+1)^2} = \int_0^1 \ln \Gamma_k(x+1) \, dx \,, \tag{4}$$

where  $\Gamma_k$  is the Bendersky generalized gamma function [3]. This function verifies in particular

$$\Gamma_k(n+1) = 1^{1^k} 2^{2^k} \cdots n^{n^k}$$
 for any integer  $n \ge 1$ .

Apparently unaware of Bendersky's work, Kurokawa and Ochiai [9, Theorem 2] have given a deep expression of the function  $\Gamma_k$  in terms of the derivative of the Hurwitz zeta function  $\zeta(s, x)$  at s = -k. Precisely, they showed that

$$\Gamma_k(x) = \exp\left\{\zeta'(-k, x) - \zeta'(-k)\right\} \quad \text{for } x > 0.$$

**Example 2.** For the first values of k, formula (4) translates into the identities

$$\sum_{n\geq 1}^{\mathcal{R}} \ln n = \ln(\sqrt{2\pi}) - 1 = \int_0^1 \ln \Gamma(x+1) \, dx \,,$$
$$\sum_{n\geq 1}^{\mathcal{R}} n \ln n = \ln(A_1) - \frac{1}{3} = \int_0^1 \ln K(x+1) \, dx$$

where  $K = \Gamma_1$  is the classical Kinkelin hyperfactorial function.

### **3** Blagouchine's first integral

We give a direct proof of [2, Theorem 2] using Cauchy's residue theorem.

**Proposition 1.** For any integer  $k \ge 0$ , let  $\mathcal{I}_k$  be the integral defined by

$$\mathcal{I}_k := \int_{-\infty}^{+\infty} \frac{\zeta(\frac{3}{2} + ix)}{(2k + 1 + 2ix)\cosh(\pi x)} \, dx \, .$$

Then

$$\mathcal{I}_k = \mu_k \tag{5}$$

where  $\mu_k$  is the conditionally convergent series defined by

$$\mu_k := \sum_{n=1}^{\infty} (-1)^{n-1} \frac{\zeta(n+1)}{n+k}$$

*Proof.* For  $k \ge 0$ , let us consider the function

$$f_k(z) = \frac{\zeta(\frac{3}{2} + iz)}{(\frac{1}{2} + k + iz)\cosh(\pi z)}.$$

We have  $\cosh(\pi z) = 0$  if and only if z = i/2 + in with  $n \in \mathbb{Z}$ . For  $n \ge 1$ , the residue of  $f_k$  at z = i/2 - in is

$$\frac{\zeta(1+n)}{(n+k)\pi\sinh(i\pi(\frac{1}{2}-n))} = \frac{\zeta(1+n)}{(n+k)i\pi\sin(\pi(\frac{1}{2}-n))} = \frac{(-1)^n\zeta(1+n)}{(n+k)i\pi}.$$

We integrate on a closed contour composed of the interval  $D_R = [-R, R]$  and the lower semicircle  $C_R$  of radius R with center at 0. By the residue theorem, we can then write the following relation:

$$\frac{1}{2i\pi} \int_{C_R} f_k(z) \, dz + \frac{1}{2i\pi} \int_{D_R} f_k(z) \, dz = -\sum_{n=1}^{N_R} \operatorname{Res}(f_k; \frac{i}{2} - in) \, ,$$

which, from the foregoing, translates into the identity

$$\int_{C_R} f_k(z) \, dz + \int_{D_R} f_k(z) \, dz = 2 \sum_{n=1}^{N_R} (-1)^{n+1} \frac{\zeta(1+n)}{(n+k)} \,. \tag{6}$$

For  $z \in C_R$ , the parameterization  $iz = Re^{it}$  with  $-\pi/2 < t < \pi/2$ , enables us to write

$$\left| \int_{C_R} f_k(z) dz \right| = \left| \int_{-\pi/2}^{+\pi/2} \frac{\zeta(\frac{3}{2} + Re^{it})}{(\frac{1}{2} + k + Re^{it}) \cosh(i\pi Re^{it})} Re^{it} dt \right|$$
$$\leq \int_{-\pi/2}^{+\pi/2} \left| \frac{\zeta(\frac{3}{2} + Re^{it})}{(\frac{1}{2} + k + Re^{it}) \cosh(i\pi Re^{it})} \right| R dt.$$

Since  $\frac{3}{2} + Re^{it}$  is in the half-plane  $\operatorname{Re}(z) > 3/2$ , its absolute value is bounded by  $\zeta(\frac{3}{2})$ , i.e.

$$\left|\zeta(\frac{3}{2} + Re^{it})\right| \le \zeta(\frac{3}{2}).$$

Hence, when R increases towards infinity, we have the following limits:

$$\lim_{R \to \infty} \int_{C_R} f_k(z) \, dz = 0 \,,$$
$$\lim_{R \to \infty} \int_{D_R} f_k(z) \, dz = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{3}{2} + ix)}{(\frac{1}{2} + k + ix) \cosh(\pi x)} \, dx = 2 \,\mathcal{I}_k \,,$$

and

$$\lim_{R \to \infty} \sum_{n=1}^{N_R} (-1)^{n+1} \frac{\zeta(1+n)}{n+k} = \sum_{n=1}^{\infty} (-1)^{n+1} \frac{\zeta(n+1)}{n+k} = \mu_k$$

This allows us to deduce formula (5) by passing to the limit in (6).

Remark 3. The constant

$$\mu_0 = \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n-1} = \sum_{n=1}^{\infty} \frac{1}{n} \ln\left(1 + \frac{1}{n}\right) = 1.257746\dots$$

has been thoroughly studied by Boyadzhiev [4] (see also [7, p. 142]). This constant is noted M in [4], K in [7], and also appears in [8] as  $\nu_{-1}$ . By a well-known series representation of Euler's constant  $\gamma$ , we also have

$$\mu_1 = \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n} = \sum_{n=1}^{\infty} \left(\frac{1}{n} - \ln\left(1 + \frac{1}{n}\right)\right) = \gamma = 0.577215\dots$$

**Example 3.** For k = 0, formula (5) translates into

$$\mathcal{I}_0 = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{3}{2} + ix)}{(1+2ix)\cosh(\pi x)} \, dx = \int_0^1 \frac{\psi(x+1) + \gamma}{x} \, dx \,,$$

where  $\psi$  is the digamma function, and for k = 1

$$\mathcal{I}_{1} = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{3}{2} + ix)}{(3 + 2ix)\cosh(\pi x)} \, dx = \gamma = -\psi(1) \, .$$

#### 4 Blagouchine's second integral

**Proposition 2.** For any integer  $k \ge 0$ , let  $\mathcal{J}_k$  be the integral defined by

$$\mathcal{J}_k := \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(2k+1+2ix)\cosh(\pi x)} \, dx$$

Then

$$\mathcal{J}_{k} = \frac{\gamma}{k+1} - \frac{1}{(k+1)^{2}} - \mathcal{I}_{k+1} \qquad (k \ge 0).$$
(7)

Since  $\mathcal{I}_1 = \gamma$ , it follows from this relation that  $\mathcal{J}_0 = -1$ ; furthermore we have the identity

$$\mathcal{J}_k = \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \ln(A_j) - \frac{k^2 + 3k + 1}{k(k+1)^2} \qquad (k \ge 1).$$
(8)

*Proof.* Formula (7) is a reformulation of [2, Theorem 1]. Now assume that  $k \ge 1$  and let  $\nu_k := \mu_{k+1}$ . By [8, Proposition 1] and (5), we have

$$\mathcal{I}_{k+1} = \nu_k = \frac{\gamma}{k+1} + \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + \frac{1}{k} + \sum_{j=0}^{k-1} \binom{k}{j} \frac{B_{j+1} H_j}{j+1}.$$

Now, using Adamchik's formula (1), this expression may be rewritten as follows:

$$\mathcal{I}_{k+1} = \frac{\gamma}{k+1} + \frac{1}{k} - \sum_{j=0}^{k-1} (-1)^j \binom{k}{j} \ln(A_j) \qquad (k \ge 1) \,.$$

Hence, formula (8) results from (7).

**Example 4.** For k = 1, formula (8) simply reduces to

$$\mathcal{J}_1 = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(3 + 2ix)\cosh(\pi x)} \, dx = \frac{1}{2}\ln(2\pi) - \frac{5}{4} \, ,$$

and for k = 2,

$$\mathcal{J}_2 = \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(5 + 2ix)\cosh(\pi x)} \, dx = \frac{1}{3}\ln(2\pi) - \frac{1}{6}\gamma + \frac{\zeta'(2)}{\pi^2} - \frac{11}{19} \, .$$

Remark 4. The following general expression can be easily deduced from formulas (2) and (3):

$$\sum_{j=0}^{k-1} (-1)^{j} {k \choose j} \ln(A_{j}) = \frac{\ln(2\pi)}{k+1} - \frac{k-1}{k+1} \frac{\gamma}{2} - \sum_{j=1}^{\left\lfloor \frac{k}{2} \right\rfloor} (-1)^{j} {k \choose 2j-1} \frac{(2j)!}{j(2\pi)^{2j}} \zeta'(2j) - \sum_{j=1}^{\left\lfloor \frac{k-1}{2} \right\rfloor} (-1)^{j} {k \choose 2j} \frac{(2j)!}{2(2\pi)^{2j}} \zeta(2j+1) \qquad (k \ge 3).$$
(9)

### 5 Further generalization

Using Fourier transform method, Candelpergher [6, Eq. (7)] recently proved the following beautiful relation which generalizes (7):

**Theorem 1.** for  $k \ge 0$  and  $\operatorname{Re}(s) > \frac{1}{2}$ , we have

$$2^{s-1}\mathcal{J}_k(s) = \frac{\gamma}{(k+1)^s} - \frac{s}{(k+1)^{s+1}} - \nu_k(s), \qquad (10)$$

with

$$\mathcal{J}_k(s) := \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(2k+1+2ix)^s \cosh(\pi x)} \, dx$$

and

$$\nu_k(s) := \sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{(n+k)^s}.$$

Applying (10) with k = 0 allows us to deduce the following identity:

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n^s} = \gamma - s - \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(\frac{1}{2} + ix)^s \cosh(\pi x)} \, dx \qquad (\operatorname{Re}(s) > \frac{1}{2}) \,. \tag{11}$$

Example 5.

$$\sum_{n=2}^{\infty} (-1)^n \frac{\zeta(n)}{n^2} = \gamma - 2 - \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(\frac{1}{2} + ix)}{(\frac{1}{2} + ix)^2 \cosh(\pi x)} \, dx \, .$$

### References

- V. Adamchik, Polygamma functions of negative order, J. Comput. Appl. Math. 100 (1998), 191–199.
- [2] I. V. Blagouchine, A complement to a recent paper on some infinite sums with the zeta values, preprint, 2020. Available at https://arxiv.org/abs/2001.00108
- [3] L. Bendersky, Sur la function gamma généralisée, Acta Math. 61 (1933), 263– 322.
- [4] K. N. Boyadzhiev, A special constant and series with zeta values and harmonic numbers, Gazeta Matematica 115 (2018), 1–16.
- [5] B. Candelpergher, Ramanujan Summation of Divergent Series, Lecture Notes in Math. 2185, Springer, 2017.
- [6] B. Candelpergher, An expansion of the Riemann Zeta function on the critical line, preprint, 2021. Available at https://hal.archives-ouvertes.fr/hal-03271709
- [7] H. Cohen, Number Theory, Volume II: Analytic and Modern Tools, Graduate Texts in Math., vol. 240, Springer, 2007.
- [8] M-A. Coppo, A note on some alternating series involving zeta and multiple zeta values, J. Math. Anal. Appl. 475 (2019), 1831–1841.
- [9] N. Kurokawa and H. Ochiai, Generalized Kinkelin's formulas Kodai Math. J. 30 (2007), 195–212.
- [10] M. Perkins and R. A. Van Gorder, Closed-form calculation of infinite products of Glaisher-type related to Dirichlet series, *Ramanujan J.* 49 (2019), 371–389.
- [11] W. Wang, Some asymptotic expansions of hyperfactorial functions and generalized Glaisher-Kinkelin constants, Ramanujan J. 43 (2017), 513–533.