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# On certain alternating series involving zeta values

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#### Introduction

This article is primarily devoted to the alternating series  $\nu_k$  defined by

$$\nu_k := \sum_{j=2}^{\infty} (-1)^j \frac{\zeta(j)}{j+k} \,,$$

where k denotes a complex parameter. By a classical result (cf. e.g. [Er], p. 45, Eq. (3) or [Jo], p. 62) which goes back to Euler's early works on zeta, one knows that  $\nu_0 = \gamma$ , where  $\gamma$  denotes the Euler-Mascheroni constant. It is less famous but yet fairly well-known (cf. [Sr], p. 135, Eq. (5.1), [SV], Eq. (1.5), [Ca], p. 93) that

$$\nu_1 = \frac{\gamma}{2} - \frac{1}{2} \ln 2\pi + 1.$$

In a recent paper, Blagouchine ([B1]) has obtained the following general expression for  $\nu_k$  in the case where k is a positive integer:

$$\nu_{k} = \frac{\gamma}{2} - \frac{\ln 2\pi}{k+1} + \frac{1}{k} + \sum_{r=1}^{\left[\frac{k}{2}\right]} (-1)^{r} {k \choose 2r-1} \frac{(2r)!}{r(2\pi)^{2r}} \zeta'(2r) + \sum_{r=1}^{\left[\frac{k+1}{2}\right]-1} (-1)^{r} {k \choose 2r} \frac{(2r)!}{2(2\pi)^{2r}} \zeta(2r+1).$$
 (1)

However, this formula is quite cumbersome. Using the functional equation of zeta, we give an equivalent but simpler expression for this series in this case. We

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also examine the series  $\nu_k$  for complex values of the parameter k, a study which, as far as we know, have never been made before. Furthermore, we study certain natural generalizations of these series involving multiple zeta values.

In the case where k is a positive integer, we show that the series  $\nu_k$  admits the following explicit evaluation:

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{2} \ln 2\pi + \sum_{j=1}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k, \qquad (2)$$

where  $C_k$  is a rational constant for which we give two different, even though equivalent, expressions (cf. Proposition 1).

A further interesting generalization of the series  $\nu_k$  may be defined as follows: for any natural number p, we consider the series

$$\nu_{k,p} := \sum_{j=2}^{\infty} \frac{(-1)^j}{j+k} \zeta(j, \underbrace{1, \dots, 1}_{p}),$$

where

$$\zeta(s_1, s_2, \cdots, s_k) = \sum_{n_1 > n_2 > \cdots > n_k \ge 1} \frac{1}{n_1^{s_1} n_2^{s_2} \cdots n_k^{s_k}},$$

in such a way that  $\nu_k = \nu_{k,0}$ . Then, for any integer  $k \ge -1$ , we show (cf. Proposition 2) the following identity:

$$\nu_{k,p} = \sum_{n=1}^{\infty} \frac{|G_n^{(k+1)}|}{n^{p+1}}.$$
 (3)

Here, the numbers  $G_n^{(k)}$  are the *Gregory coefficients of higher order* introduced in [B2]. They may be either defined by

$$G_n^{(k)} := \frac{1}{n!} \sum_{m=1}^n \frac{S_1(n,m)}{m+k}, \quad \text{for } k = 0, 1, 2, \dots$$
 (4)

where  $S_1(n,m)$  are the Stirling numbers of the first kind, or equivalently by the integral formula

$$G_n^{(k)} = \frac{(-1)^n}{n!} \int_0^1 x^{k-1} (-x)_n \, dx \,, \tag{5}$$

where  $(z)_n = z(z+1)(z+2)\cdots(z+n-1)$  is the Pochhammer symbol. It follows from this last expression that  $G_n^{(k)} = (-1)^{n+1}|G_n^{(k)}|$ . Furthermore,  $G_n^{(1)} = G_n$ , where  $G_n$  are the *Gregory coefficients*, also known as the *Bernoulli numbers of the second kind*  $^1$ .

<sup>&</sup>lt;sup>1</sup>Several authors quoted in reference use different notations for these numbers, they are noted  $b_n$  in [Jo], [CY], and  $\beta_n/n!$  in [Ca].

As a special case of the identity above, we deduce that

$$\nu_{k-1} = \sum_{n=1}^{\infty} \frac{|G_n^{(k)}|}{n}, \quad \text{for } k = 0, 1, 2, \dots$$
 (6)

and we note that this formula nicely generalizes the famous Mascheroni series for  $\gamma$  (which is nothing else than the case k = 1).

Finally, the extension to the complex case is studied in the last section. We provide two beautiful integral formulas for the same series:

$$\nu_k = \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(3/2 \pm ix)}{(3/2 \pm ix + k) \cosh \pi x} dx, \qquad (7)$$

and

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{(k+1)^2} - \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(1/2 \pm ix)}{(1/2 \pm ix + k) \cosh \pi x} dx, \qquad (8)$$

which are valid for  $\text{Re}(k) > -\frac{3}{2}$  and  $\text{Re}(k) > -\frac{1}{2}$  respectively. The second representation seems especially interesting because the integral runs over the whole critical line.

In appendix, we return to the case of a positive integer k and highlight the existence of an interesting relation between the series  $\nu_k$ , the Stirling numbers of the second kind  $S_2(n, k)$  and the shifted Mascheroni series  $\sigma_r$  which have been recently studied in [CY].

## 1 The case of a positive integer

In this section, we study the case where the parameter k is a positive integer and give two independant proofs of our formula (2). More precisely, we prove the following proposition:

**Proposition 1.** For any positive integer k, then

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{2} \ln 2\pi + \sum_{i=1}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k$$

with

$$C_{k} = \frac{1}{k} + \sum_{r=1}^{\left[\frac{k}{2}\right]} {k \choose 2r} \frac{B_{2r} H_{2r-1}}{k+1-2r}$$

$$= \frac{H_{k}}{k+1} + \frac{1}{2k} - \sum_{r=1}^{\left[\frac{k}{2}\right]} \frac{B_{2r}}{2r(k+1-2r)}$$
(9)

where  $H_n$  are the harmonic numbers defined by

$$H_n = \sum_{m=1}^n \frac{1}{m}$$
, for  $n = 1, 2, 3, \dots$ 

and  $B_n$  are the Bernoulli numbers defined by their exponential generating series

$$\sum_{n=0}^{\infty} B_n \frac{z^n}{n!} = \frac{z}{e^z - 1} \,, \quad \text{for } |z| < 2\pi \,.$$

In particular,  $B_0 = 1$ ,  $B_1 = -1/2$ ,  $B_{2r+1} = 0$  for  $r \ge 1$ .

*Proof.* We can quite easily deduce (2) from formula (1). A differentiation of the functional equation

$$\zeta(s) = 2(2\pi)^{s-1}\Gamma(1-s)\zeta(1-s)\sin\frac{\pi s}{2},$$

leads to the relations

$$(-1)^r \frac{(2r)!}{2(2\pi)^{2r}} \zeta(2r+1) = \zeta'(-2r), \text{ for } r=1,2,3,\cdots$$

and

$$(-1)^r \frac{(2r)!}{r(2\pi)^{2r}} \zeta'(2r) = -\zeta'(1-2r) + \frac{B_{2r}}{2r} \left( H_{2r-1} - \gamma - \ln 2\pi \right), \quad \text{for } r = 1, 2, 3, \cdots.$$

Substituting these relations in (1) and aggregating the terms under the two  $\Sigma$  then gives

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{2} \ln 2\pi + \sum_{j=1}^{k-1} (-1)^j \binom{k}{j} \zeta'(-j) + C_k$$

with

$$C_k = \frac{1}{k} + \sum_{r=1}^{\left[\frac{k}{2}\right]} {k \choose 2r} \frac{B_{2r} H_{2r-1}}{k+1-2r}.$$

Another proof of formula (2), independent from (1), may also be deduced from the following development given in [Ca] p. 93:

$$\sum_{k=0}^{\infty} \frac{(-1)^k z^k}{k!} \sum_{j=1}^{\infty} \frac{(-1)^{j-1}}{j} \zeta^{\mathcal{R}}(j-k) = (1-e^z) \sum_{k=0}^{\infty} \frac{(-1)^k z^k}{k!} \zeta'(-k) + (1-e^z) \sum_{k=0}^{\infty} \frac{(-1)^k z^k}{k!} \frac{1}{(k+1)^2} + \int_0^1 \ln(t+1)e^{-zt} dt,$$

with

$$\zeta^{\mathcal{R}}(j-k) = \begin{cases} \gamma & \text{if } j = k+1\\ \zeta(j-k) - \frac{1}{j-k-1} & \text{otherwise.} \end{cases}$$

Rewriting the series  $\nu_k$  under the following form:

$$\nu_k = \sum_{j=k+2}^{\infty} \frac{(-1)^{j-k}}{j} \zeta(j-k),$$

and using the well-known relations:

$$\zeta(0) = -\frac{1}{2}, \quad \zeta'(0) = -\frac{1}{2} \ln 2\pi, \quad \zeta(1 - 2r) = -\frac{B_{2r}}{2r},$$

as well as the combinatorial identity

$$\sum_{j=0}^{k} (-1)^j \binom{k}{j} \frac{1}{(j+1)^2} = \frac{H_{k+1}}{k+1} ,$$

then, a careful identification of the terms in  $\frac{z^k}{k!}$  in the previous development leads to the same formula (2) with a simpler expression of the constant  $C_k$ :

$$C_k = \frac{H_k}{k+1} + \frac{1}{2k} - \sum_{r=1}^{\left[\frac{k}{2}\right]} \frac{B_{2r}}{2r(k+1-2r)}.$$

**Example 1.** For the first six values of k, we obtain the following relations:

$$\begin{split} \nu_1 &= \frac{\gamma}{2} - \frac{1}{2} \ln 2\pi + 1 \,, \\ \nu_2 &= \frac{\gamma}{3} - \frac{1}{2} \ln 2\pi - 2\zeta'(-1) + \frac{2}{3} \,, \\ \nu_3 &= \frac{\gamma}{4} - \frac{1}{2} \ln 2\pi - 3\zeta'(-1) + 3\zeta'(-2) + \frac{7}{12} \,, \\ \nu_4 &= \frac{\gamma}{5} - \frac{1}{2} \ln 2\pi - 4\zeta'(-1) + 6\zeta'(-2) - 4\zeta'(-3) + \frac{47}{90} \,, \\ \nu_5 &= \frac{\gamma}{6} - \frac{1}{2} \ln 2\pi - 5\zeta'(-1) + 10\zeta'(-2) - 10\zeta'(-3) + 5\zeta'(-4) + \frac{167}{360} \,, \\ \nu_6 &= \frac{\gamma}{7} - \frac{1}{2} \ln 2\pi - 6\zeta'(-1) + 15\zeta'(-2) - 20\zeta'(-3) + 15\zeta'(-4) - 6\zeta'(-5) + \frac{349}{840} \,. \end{split}$$

# 2 Alternating series involving multiple zeta values

In this section, we generalize the series  $\nu_k$  to certain multiple zeta values and prove our formulae (3) and (6).

**Proposition 2.** For all natural numbers  $p \geq 0$  and integers  $k \geq -1$ , then

$$\nu_{k,p} := \sum_{j=2}^{\infty} \frac{(-1)^j}{j+k} \zeta(j, \underbrace{1, \dots, 1}_{p}) = \sum_{n=1}^{\infty} \frac{|G_n^{(k+1)}|}{n^{p+1}},$$

where  $G_n^{(k)}$  are the Gregory coefficients of higher order defined by Eq. (5) or (6). In particular, for  $k \geq 0$ ,

$$\nu_{k-1} = \sum_{n=1}^{\infty} \frac{|G_n^{(k)}|}{n} \,.$$

**Example 2.** For the first values of k, we obtain the following developments in series containing only positive rational terms:

$$\begin{split} \nu_{-1} &= 1 + \frac{1}{8} + \frac{5}{108} + \frac{3}{128} + \frac{251}{18000} + \frac{95}{10368} + \dots, \\ \nu_{0} &= \frac{1}{2} + \frac{1}{24} + \frac{1}{72} + \frac{19}{2880} + \frac{3}{800} + \frac{863}{362880} + \dots, \\ \nu_{1} &= \frac{1}{3} + \frac{1}{48} + \frac{7}{1080} + \frac{17}{5760} + \frac{41}{25200} + \frac{731}{725760} + \dots, \\ \nu_{2} &= \frac{1}{4} + \frac{1}{80} + \frac{1}{270} + \frac{11}{6720} + \frac{89}{100800} + \frac{5849}{10886400} + \dots, \\ \nu_{3} &= \frac{1}{5} + \frac{1}{120} + \frac{1}{420} + \frac{83}{80640} + \frac{59}{108000} + \frac{397}{1209600} + \dots, \\ \nu_{4} &= \frac{1}{6} + \frac{1}{168} + \frac{5}{3024} + \frac{17}{24192} + \frac{557}{1512000} + \frac{5249}{23950080} + \dots, \\ \nu_{5} &= \frac{1}{7} + \frac{1}{224} + \frac{11}{9072} + \frac{41}{80640} + \frac{439}{1663200} + \frac{311}{1995840} + \dots, \\ \nu_{6} &= \frac{1}{8} + \frac{1}{288} + \frac{1}{1080} + \frac{73}{190080} + \frac{47}{237600} + \frac{2581}{22239360} + \dots, \end{split}$$

*Proof.* In order to prove Proposition 2 above, we use the following lemma (cf. [Xu] Eq. (2.27) and (2.28)):

**Lemma 1.** For all integers  $m \ge 1$  and  $p \ge 0$ , one has

$$\int_0^1 \frac{\ln^m (1-x) \ln^p (x)}{x} dx = (-1)^{m+p} m! \, p! \, \zeta(m+1, \underbrace{1, \dots, 1}_{p}).$$

Then, we can write the following equalities:

$$\begin{split} &\sum_{j=2}^{\infty} \frac{(-1)^j}{j+k} \zeta(j, \underbrace{1, \dots, 1}) = \sum_{m=1}^{\infty} \frac{(-1)^{m+1}}{m+k+1} \zeta(m+1, \underbrace{1, \dots, 1}) \\ &= \frac{(-1)^{p+1}}{p!} \sum_{m=1}^{\infty} \frac{1}{m+k+1} \int_0^1 \frac{\ln^m (1-x)}{m!} \frac{\ln^p (x)}{x} \, dx \\ &= \frac{(-1)^{p+1}}{p!} \sum_{m=1}^{\infty} \frac{(-1)^m}{m+k+1} \int_0^1 \left( \sum_{n=1}^{\infty} |S_1(n,m)| \frac{x^n}{n!} \right) \frac{\ln^p (x)}{x} \, dx \\ &= -\sum_{n=1}^{\infty} \frac{1}{n!} \sum_{m=1}^n \frac{(-1)^m}{m+k+1} |S_1(n,m)| \frac{(-1)^p}{p!} \int_0^1 x^{n-1} \ln^p (x) \, dx \\ &= \sum_{n=1}^{\infty} (-1)^{n+1} \left( \frac{1}{n!} \sum_{m=1}^n \frac{S_1(n,m)}{m+k+1} \right) \frac{1}{n^{p+1}} \\ &= \sum_{n=1}^{\infty} (-1)^{n+1} \frac{G_n^{(k+1)}}{n^{p+1}} = \sum_{n=1}^{\infty} \frac{|G_n^{(k+1)}|}{n^{p+1}} \, . \end{split}$$

Example 3.

$$\sum_{j=2}^{\infty} \frac{(-1)^j}{j} \sum_{n=1}^{\infty} \frac{H_n}{n^j} = \nu_{0,1} - \nu_{-1} + \zeta(2) = \sum_{n=1}^{\infty} \frac{|G_n|}{n^2} - \sum_{n=1}^{\infty} \frac{|G_n^{(0)}|}{n} + \frac{\pi^2}{6}.$$

## 3 Integral representations

In this section, we study the extension to the case of a complex parameter k and give two beautiful integral representations for the series  $\nu_k$ .

**Proposition 3.** For any  $k \in \mathbb{C}$  such that Re(k) > -3/2, then

$$\nu_k = \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(3/2 \pm ix)}{(3/2 \pm ix + k) \cosh \pi x} dx.$$
 (10)

For any  $k \in \mathbb{C}$  such that Re(k) > -1/2, then

$$\nu_k = \frac{\gamma}{k+1} - \frac{1}{(k+1)^2} - \frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(1/2 \pm ix)}{(1/2 \pm ix + k) \cosh \pi x} dx \tag{11}$$

By identification with formula (2), we deduce the following corollary:

Corollary 1. For any positive integer k, we have

$$\frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(1/2 \pm ix)}{(1/2 \pm ix + k) \cosh \pi x} dx = \sum_{j=0}^{k-1} (-1)^{j+1} {k \choose j} \zeta'(-j) - C_k - \frac{1}{(k+1)^2},$$

where  $C_k$  is the rational constant defined by (9).

#### Example 4.

$$\frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(3/2 \pm ix)}{(3/2 \pm ix) \cosh \pi x} dx = \gamma,$$

$$\frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(3/2 \pm ix)}{(1/2 \pm ix) \cosh \pi x} dx = \nu_{-1},$$

$$\frac{1}{2} \int_{-\infty}^{+\infty} \frac{\zeta(1/2 \pm ix)}{(1/2 \pm ix) \cosh \pi x} dx = -1,$$

$$\frac{1}{2} \int_{-\infty}^{\infty} \frac{\zeta(1/2 \pm ix)}{(3/2 + ix) \cosh(\pi x)} dx = \frac{1}{2} \ln 2\pi - \frac{5}{4},$$

$$\frac{1}{2} \int_{-\infty}^{\infty} \frac{\zeta(1/2 + ix)}{(5/2 + ix) \cosh(\pi x)} dx = \frac{1}{2} \ln 2\pi - 2 \ln(A) - \frac{11}{18},$$

where  $A := \exp\left\{\frac{1}{12} - \zeta'(-1)\right\}$  is the Glaisher-Kinkelin constant.

*Proof.* In order to prove Proposition 3, we resort to the contour integration method. Consider the following line integral taken along a contour C consisting of the interval [-R, +R],  $R \in \mathbb{N}$ , on the real axis, and a semicircle of the radius R in the upper half-plane, denoted  $C_R$ ,

$$\oint_{C} \frac{\zeta(1/2 - iz)}{(1/2 - iz + k) \cosh \pi z} dz = \int_{-R}^{+R} \frac{\zeta(1/2 - ix)}{(1/2 - ix + k) \cosh \pi x} dx + \int_{C_{R}} \frac{\zeta(1/2 - iz)}{(1/2 - iz + k) \cosh \pi z} dz.$$
(12)

On the contour  $C_R$  the last integral may be bounded as follows:

$$\left| \int_{C_R} \frac{\zeta(1/2 - iz)}{(1/2 - iz + k) \cosh \pi z} dz \right| = R \left| \int_0^{\pi} \frac{\zeta(1/2 - iRe^{i\varphi}) e^{i\varphi}}{(1/2 - iRe^{i\varphi} + k) \cosh (\pi Re^{i\varphi})} d\varphi \right| \le$$

$$\le R \max_{\varphi \in [0, \pi]} \left| \frac{\zeta(1/2 - iRe^{i\varphi})}{1/2 - iRe^{i\varphi} + k} \right| \cdot I_R \le \max_{\varphi \in [0, \pi]} \left| \zeta(1/2 - iRe^{i\varphi}) \right| \cdot I_R$$
(13)

where we denoted

$$I_R := \int_0^{\pi} \frac{d\varphi}{\left|\cosh\left(\pi R e^{i\varphi}\right)\right|}, \qquad R > 0,$$

for the purpose of brevity. Now, in the half-plane  $\sigma > 1$ , the absolute value of  $\zeta(\sigma+it)$  may be always bounded by a constant  $C=\zeta(\sigma)$ , which decreases and tends to 1 as  $\sigma \to \infty$ . In contrast, in the strip  $0 \leqslant \sigma \leqslant 1$  the upper bound of the function  $\left|\zeta(\sigma+it)\right|$  grows, and presently it is not known which is its exact rate of grow. However, it follows from the general theory of Dirichlet series that it cannot be faster than  $O(t^{1/4})$ . Hence, since  $\sin \varphi \geqslant 0$  and if R is large enough, this rough estimate gives us

$$\left|\zeta\left(\frac{1}{2}-iRe^{i\varphi}\right)\right| = \left|\zeta\left(\frac{1}{2}+R\sin\varphi-iR\cos\varphi\right)\right| = O\left(\sqrt[4]{R}\right)$$

in the interval  $\varphi \in [0, \pi]$ . On the other hand, as R tends to infinity and remains integer the integral  $I_R$  tends to zero as O(1/R). To show this, it is sufficient to remark that

$$\frac{1}{\left|\cosh\left(\pi R e^{i\varphi}\right)\right|} = \frac{\sqrt{2}}{\sqrt{\cosh(2\pi R\cos\varphi) + \cos(2\pi R\sin\varphi)}} = O\left(e^{-\pi R|\cos\varphi|}\right), \qquad R \to \infty,$$

since  $0 \leqslant \varphi \leqslant \pi$  and R is integer. Thus, accounting for the symmetry of  $\left|\cosh\left(\pi Re^{i\varphi}\right)\right|^{-1}$  about  $\varphi = \frac{1}{2}\pi$ , we deduce that

$$I_{R} = \int_{0}^{\pi} \frac{\sqrt{2}}{\sqrt{\cosh(2\pi R \cos \varphi) + \cos(2\pi R \sin \varphi)}} d\varphi$$

$$= \int_{0}^{\frac{\pi}{2}} \frac{2\sqrt{2}}{\sqrt{\cosh(2\pi R \cos \varphi) + \cos(2\pi R \sin \varphi)}} d\varphi$$

$$= O\left(\int_{0}^{\frac{\pi}{2}} e^{-\pi R \cos \varphi} d\varphi\right) = O\left(\int_{0}^{\frac{\pi}{2}} e^{-\pi R \sin \vartheta} d\vartheta\right), \qquad R \to \infty.$$
 (14)

From the well–known inequality

$$\frac{2\vartheta}{\pi} \leqslant \sin\vartheta \leqslant \vartheta \,, \qquad \vartheta \in \left[0, \frac{1}{2}\pi\right]$$

it finally follows that

$$\frac{1 - e^{-\frac{1}{2}\pi^2 R}}{\pi R} \leqslant \int_{0}^{\frac{\pi}{2}} e^{-\pi R \sin \vartheta} \, d\vartheta \leqslant \frac{1 - e^{-\pi R}}{2R} \,, \tag{15}$$

and since R is large, exponential terms on both sides may be neglected. Thus  $I_R = O(1/R)$  at  $R \to \infty$ .<sup>2</sup> Inserting both latter results into (13), we obtain

$$\left| \int_{C_R} \frac{\zeta(1/2 - iz)}{(1/2 - iz + k) \cosh \pi z} dz \right| = O\left(R^{-3/4}\right) \to 0 \quad \text{as} \quad R \to \infty, R \in \mathbb{N},$$

Hence, making  $R \to \infty$ , equality (12) becomes

$$\int_{-\infty}^{+\infty} \frac{\left(a - ix\right)^{1-s}}{\cosh^2 \pi x} dx = \oint_C \frac{\left(a - iz\right)^{1-s}}{\cosh^2 \pi z} dz \tag{17}$$

where the latter integral is taken around an infinitely large semicircle in the upper half-plane. The integrand is not a holomorphic function: it has the poles of the second order at  $z=z_n\equiv i\left(n+\frac{1}{2}\right),\ n\in\mathbb{N}_0$ , due to the hyperbolic secant, and a branch point at z=-ia due to the term in the numerator. If  $\mathrm{Re}(a)>0$ , the branch point lies outside the integration contour and we may use the Cauchy residue theorem:

$$\oint_{C} \frac{\left(a - iz\right)^{1-s}}{\cosh^{2}\pi z} dz = 2\pi i \sum_{n=0}^{\infty} \underset{z=z_{n}}{\operatorname{res}} \frac{\left(a - iz\right)^{1-s}}{\cosh^{2}\pi z} =$$

$$= -\frac{2i}{\pi} \sum_{n=0}^{\infty} \frac{\partial}{\partial z} \left(a - iz\right)^{1-s} \Big|_{z=i\left(n+\frac{1}{2}\right)} =$$

$$= \frac{2(s-1)}{\pi} \sum_{n=0}^{\infty} \left(a + \frac{1}{2} + n\right)^{-s} = \frac{2(s-1)}{\pi} \zeta(s, a + \frac{1}{2}).$$
(18)

$$\int_0^{\frac{\pi}{2}} e^{-\pi R \sin \vartheta} d\vartheta = \frac{\pi}{2} \left\{ I_0(\pi R) - L_0(\pi R) \right\} \sim \frac{1}{\pi R} , \qquad R \to \infty , \tag{16}$$

i.e. the integral asymptotically tends to the left bound (15).

<sup>&</sup>lt;sup>2</sup>Another way to obtain the same result is to recall that the integral (15) may be evaluated in terms of the modified Bessel function  $I_n(z)$  of the first kind and the modified Struve function  $I_n(z)$ . Using the asymptotic expansions of these special functions we obtain even a more exact result, namely

Equating (17) with the last result yields

$$\zeta(s, a + \frac{1}{2}) = \frac{\pi}{2(s-1)} \int_{-\infty}^{+\infty} \frac{\left(a - ix\right)^{1-s}}{\cosh^2 \pi x} dx, \quad \text{Re}(a) > 0.$$
 (19)

Splitting the interval of integration in two parts  $(-\infty, 0]$  and  $[0, +\infty]$  and recalling that

$$\left(a+ix\right)^{s} + \left(a-ix\right)^{s} = 2\left(a^{2}+x^{2}\right)^{\frac{s}{2}}\cos\left(s\operatorname{arctg}\frac{x}{a}\right) \tag{20}$$

the latter expression may also be written as

$$\zeta(s, a + \frac{1}{2}) = \frac{\pi}{2(s-1)} \int_0^\infty \frac{(a+ix)^{1-s} + (a-ix)^{1-s}}{\cosh^2 \pi x} dx$$
 (21)

$$= \frac{\pi}{s-1} \int_0^\infty \frac{\cos\left[(s-1)\arctan\frac{x}{a}\right]}{\left(a^2+x^2\right)^{\frac{1}{2}(s-1)}\cosh^2\pi x} dx, \quad \text{Re}(a) > 0. \quad (22)$$

Appendix

Let us consider the forward shifted Mascheroni series which are defined by

$$\sigma_r := \sum_{n=1}^{\infty} \frac{|G_{n+r}|}{n}, \text{ for } r = 0, 1, 2, \cdots.$$

These series were the main subject of a recent article ([CY]). We have among other things established the following decomposition of  $\zeta'(-j)$  (cf. [CY], Proposition 3):

$$\zeta'(-j) = \sum_{r=2}^{j+1} (-1)^{j-r} (r-1)! S_2(j, r-1) \sigma_r - \frac{B_{j+1}}{j+1} \gamma - \frac{B_{j+1}}{(j+1)^2}, \quad \text{for } j = 1, 2, 3, \dots$$

where  $S_2(j,r)$  are the Stirling numbers of the second kind. Then, substituting this relation in (2) enables to write the series  $\nu_k$  as an integer linear combination of  $\sigma_1, \sigma_2, \dots, \sigma_k$  and a rational constant  $D_k$  linked to  $C_k$  (the coefficient of  $\gamma = \sigma_0$  vanishes by a well-known relation between the Bernoulli numbers<sup>3</sup>). More precisely, we obtain the following nice rewriting of formula (2):

The coefficient of  $\gamma$  is  $\frac{1}{k+1} \sum_{j=0}^{k} {k+1 \choose j} B_j$  which is equal to zero.

$$\nu_k = D_k + \sum_{r=1}^k (-1)^r (r-1)! \left( \sum_{j=r-1}^{k-1} {k \choose j} S_2(j,r-1) \right) \sigma_r, \text{ for } k = 1, 2, 3, \dots$$

with

$$D_k = C_k - \frac{1}{2} + \sum_{r=1}^{\left[\frac{k}{2}\right]} {k \choose 2r} \frac{B_{2r}}{2r(k+1-2r)} = \frac{1}{k} - \frac{1}{2} + \sum_{r=1}^{\left[\frac{k}{2}\right]} {k \choose 2r} \frac{B_{2r}H_{2r}}{k+1-2r} \,. \tag{23}$$

**Example 5.** For the first six values of k, we obtain the following relations:

$$\begin{split} \nu_1 &= \frac{1}{2} - \sigma_1 \,, \\ \nu_2 &= \frac{1}{4} - \sigma_1 + 2\sigma_2 \,, \\ \nu_3 &= \frac{5}{24} - \sigma_1 + 6\sigma_2 - 6\sigma_3 \,, \\ \nu_4 &= \frac{13}{72} - \sigma_1 + 14\sigma_2 - 36\sigma_3 + 24\sigma_4 \,, \\ \nu_5 &= \frac{109}{720} - \sigma_1 + 30\sigma_2 - 150\sigma_3 + 240\sigma_4 - 120\sigma_5 \,, \\ \nu_6 &= \frac{23}{180} - \sigma_1 + 62\sigma_2 - 420\sigma_3 + 1560\sigma_4 - 1800\sigma_5 + 720\sigma_6 \,. \end{split}$$

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